Potential formulation (U, A): 1/3

In Electrostatics $\nabla x \bar{E} = 0$ allowed Us to write \bar{E} as a gradient of a Scalar protential $\bar{E} = -\bar{\nabla}U$

In electrodynamics this is no linger possible, because the $\nabla \times E$ is non-zero. But B vernains divergence less. So we can still write

as in magnetostatics putting it in to the Favaday's law field we get $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\frac{\partial}{\partial t} (\nabla \times \vec{A}).$ $\nabla \times \vec{E} + \nabla \times \partial \vec{A} = 0$ $\nabla \times (\vec{E} + \partial \vec{A}) = 0$

Here is a quantity unlike É alone, Whose curl does not vanish et can be written as the gradient of a Scalar potential

2/3

E+DAZ -FU

E = - FU - 2A -- (2)

This reduces the above equations, to old form of course in electrostatics as we put $\partial A = 0$

The potential formulation of equation (1) & (2) automatically full fills the two homogeness Maxwell's equations

F. EZ S. & TXB=HOJ+MODD for free space.

Putting equatur (2), the Evalue in Gauss's law for electrostatics

EMPF-1 we have F. (-84-2A)= 3/3 or 7 u + 2 ((√ x A) = - - - - - - - - - - - - 3) Equation (3) replaces Poisson's Equation. Let us put the relevant terms from D& 2) in Maxwell > 4th egnetim OXB = HOT+ EOMO DE Ruthy BZTXA & E=-FU-JA we get TXTXA = MOJ-Moto FOUN - 640 2A Using Veilor Identity DXDXA = D(D.A) - JA ---(4) we have VA - E.M. 2A - TI (T. A+MOE) = - HOJ - (S) Equation (3) & (5) carry all the Information in Maxwell's equations in electrodynamics à hence termed as protential firmulation of u & A.